

## Variables of State:

e.g.  $E = \text{internal energy}$   $\leftarrow$  describes a state  
also  $P, V, T, n, m(\text{mass})$

Not variables of state:

Heat =  $q$  describe a specific way to  
Work =  $w$  change  $E$ . Can get  $\xrightarrow{\text{to}}$   
from  $E_1$  to  $E_2$  through all  
heat, all work, or any  
combination.

## Define a NEW variable of state:

Enthalpy  $H = E + PV$  formed from a combination  
of state variables

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## Two classes:

Extensive  $\Rightarrow$  proportional to mass of the system  
e.g.  $V, E, H, C$

Intensive  $\Rightarrow$  independent of mass of the system  
e.g.  $P, T, \text{molar heat capacity}, \text{specific heat}$

## HEAT AND WORK — TRANSFERS OF ENERGY

$$\text{Change in energy} \Rightarrow E_2 - E_1 = q + w$$

↑                              ↑                              ↑  
 Energy of final state      energy of initial state      heat transferred  
 to the system  
 (heat flows to)              work done  
 on the system  
 (work to)

ASIDE: THIS APPLIES FOR A CLOSED SYSTEM

(energy, but not matter can be exchanged with surroundings).

THIS IS THE FIRST LAW

(and  $E_1$  and  $E_2$ )

The important point here is that  $E$  is a state variable,  $q$  and  $w$  are NOT

You can raise the energy of the system by:

- 1) transferring  $E$  as heat only
- 2) transferring  $E$  as work only
- 3) transferring some of each

Can unfold a protein by heating ( $T$ -jump)

Can unfold a protein by increasing pressure ( $P$ -jump)

Either way puts energy into the system, raising the energy of the protein

(3)

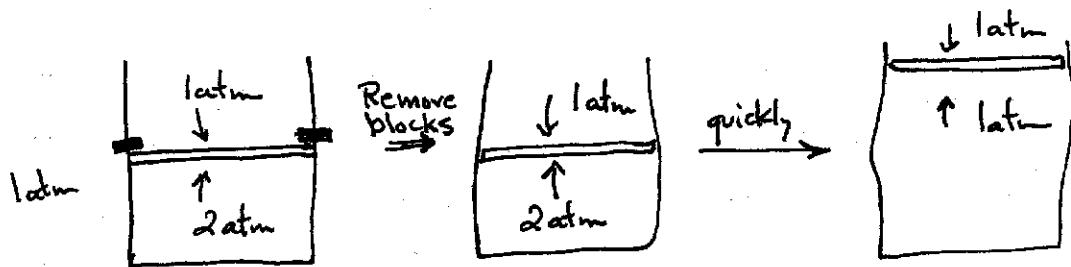
Aside: note that some texts (the minority) define work in an opposite sense, so that  $E = q - w \Leftarrow$  WE WILL NOT

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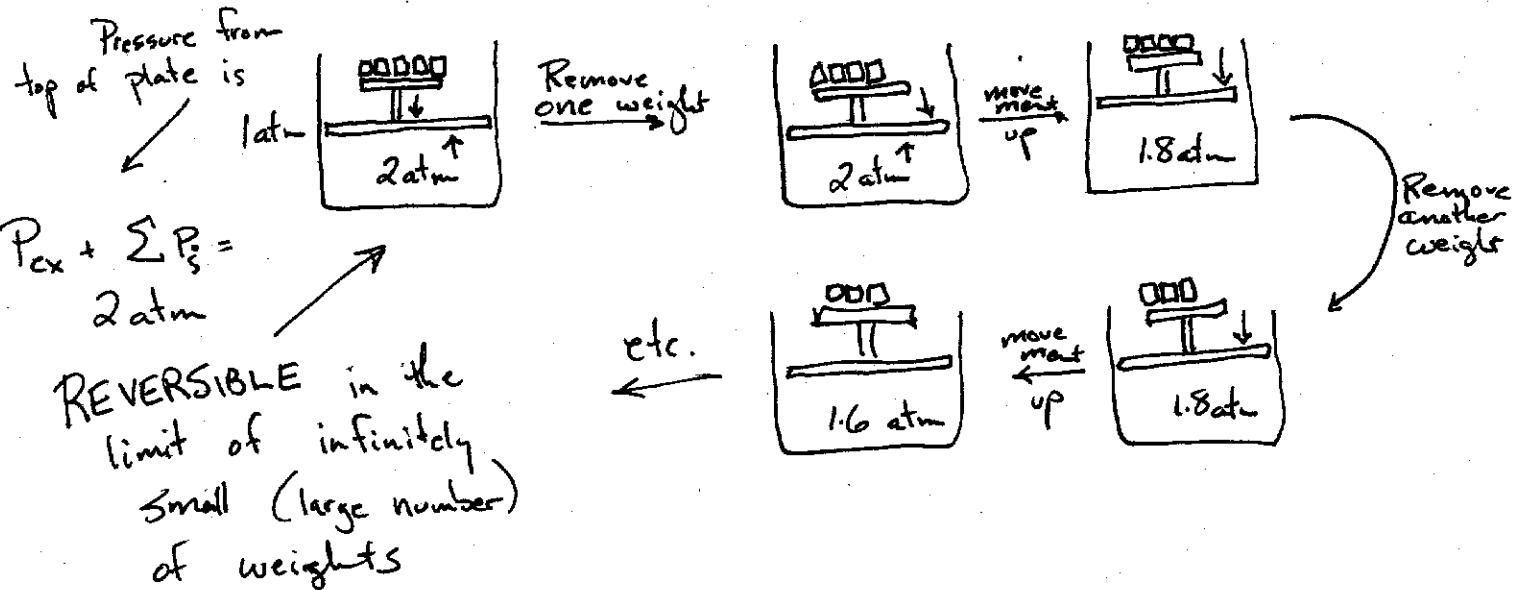
### Reversible Path

There are many ways of getting from one state (e.g.  $E_1$ ) to another (e.g.  $E_2$ ).

e.g.



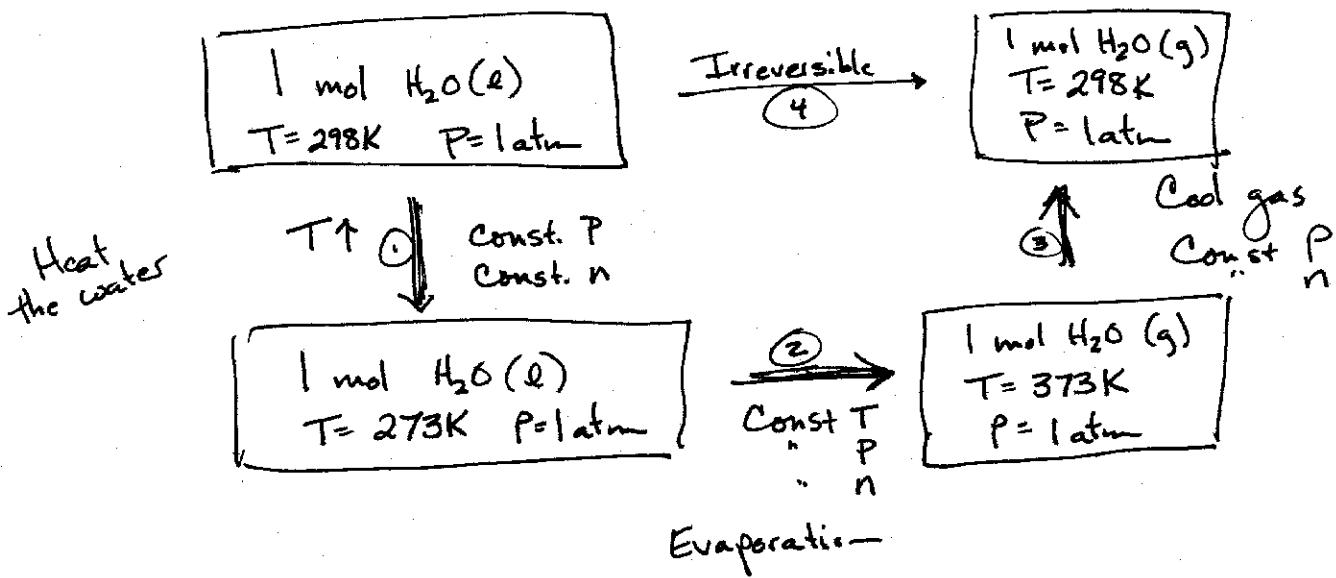
IRREVERSIBLE expansion  $\rightarrow$  after blocks are removed pressure difference drives rapid movement of the piston. Pressure on two sides of piston plate not equal.



A reversible path is one in which the process can be reversed at any point along the transition (path).

Simpler example:

At  $100^{\circ}\text{C}$  and 1 atm, the phase transition is poised between liquid and gaseous water (steam)



Constant  $P$  = isobaric

Constant  $T$  = isothermal

Constant volume = isochoric  
No heat transferred ( $q=0$ ) = adiabatic

Processes ① and ② are easily imagined.

Process ③ less so

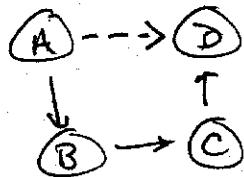
↑  
isolated system

Processes ①, ②, and ③ are reversible

Process ④ is not.

Measurement is often easier along reversible paths

Why construct these paths?



We may be interested in the energetics of  $A \rightarrow D$ , but we cannot conduct the measurement.

But if we get from  $A \rightarrow \dots \rightarrow D$  by other measurable paths, then the net energetics that we measure provide the original measure we sought.

This is only true iff we are measuring a STATE VARIABLE of the system

### Effects of Changes in T and P

Why do we care?

We know that proteins (and DNA) denature at high temperature

- 1) Cooking an egg (protein denatures)
- 2) PCR — reversible denaturation of DNA (but not the protein!)

Extremophiles  $\Rightarrow$  organisms that grow (and have stable proteins) at high T and P.

We want to know how!?

(6)

p. 33 (Table 2.2)

Overhead?

Compares properties of liquid, solid, gaseous water.

ASIDE: Water has unique properties which dictate most biological interactions.  
Life evolved in any other solvent would be very different!

Back to heat capacities →

Heat <sub>water</sub><sup>1 mol</sup> from 0°C → 100°C (l→l) [1 atm]

$$q_p = \int_{273}^{373} C_p dT$$

Constant pressure

$$C_p = (1 \text{ mol}) \times (75.4 \text{ J K}^{-1} \text{ mol}^{-1}) = 75.4 \text{ J K}^{-1}$$

≈ constant over range at 1 atm

$$\therefore q_p = (75.4 \text{ J K}^{-1}) [373 - 273] \text{ K} = 7540 \text{ J}$$

What about w<sub>p</sub>?

$$w_p = \int p dV = (1 \text{ atm}) \int_{V_1}^{V_2} dV$$

$$\text{Volume} = \frac{\text{mass}}{\text{density}} = \frac{(M_w)(\text{moles})}{\text{density}}$$

$$= (1 \text{ atm})(V_2 - V_1)$$

Look  
in table 2.2

$$\therefore V_2 - V_1 = (M_w)(\text{moles}) \left[ \frac{1}{\rho_{100}} - \frac{1}{\rho_0} \right]$$

$$= (18 \text{ g mol}^{-1})(1 \text{ mol}) \left[ \frac{1}{0.9584} - \frac{1}{0.9999} \right] \text{ g cm}^{-3}$$

$$= 0.78 \text{ cm}^3$$

$$(1.0434 - 1.0000)$$

$$w_p = (1 \text{ atm})(0.78 \text{ cm}^3) \left[ 0.1013 \frac{\text{J}}{\text{cm}^3 \cdot \text{atm}} \right]$$

$$= 0.079 \text{ J}$$

Do for homework

Contrast this with  $q_p = 7540 \text{ J}$

i.e. all the change in  $E$  is from  $q_p$

$$\therefore E_2 - E_1 = q_p + w_p = 7540 \text{ J} = \Delta E$$

Similarly,

$$H_2 - H_1 = E_2 + P_2 V_2 - E_1 - P_1 V_1$$

$$H = E + PV$$

$$= E_2 - E_1 - (P)(V_2 - V_1)$$

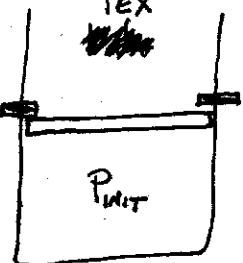
$$\Delta H = 7540 \text{ J} \quad (\text{actually, exactly})$$

Why?

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## PV Work and Gas Expansion

Irreversible



$$\omega = - \int P dV \quad P = P_{\text{ext}} = \text{const.}$$

$$= -P(V_2 - V_1)$$

Reversible, isothermal

$$w = - \int P_{\text{ex}} dV$$

Pressure slowly changes, temperature constant

From ideal gas law  $P = \frac{nRT}{V}$

$$\begin{aligned} w_T &= - \int \frac{nRT}{V} dV \quad \text{but } n, R, \text{ and } T \text{ constant} \\ &= -nRT \int_{V_1}^{V_2} \frac{dV}{V} \\ &= -nRT \ln \frac{V_2}{V_1} = -nRT \ln \frac{P_1}{P_2} \end{aligned}$$

WHY? ↪

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So work irreversible  $\Rightarrow -P(V_2 - V_1)$

work reversible  $\Rightarrow -nRT \ln \frac{V_2}{V_1}$

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For an IDEAL GAS (no inter-molecule interactions),  
E should be independent of pressure

$$\text{So } E_2 - E_1 = 0 = q_T + w_T$$

temperature constant  
pressure irrelevant

$$\therefore q_T = -w_T = +nRT \ln \frac{V_2}{V_1}$$